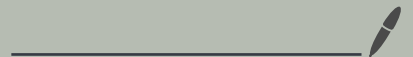


Analysis IV

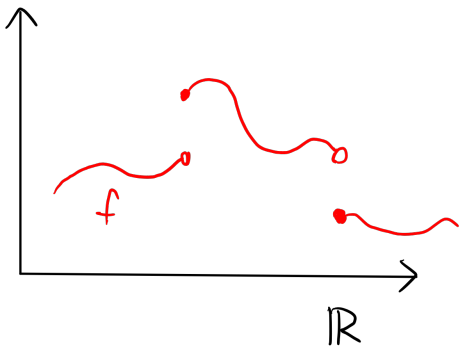
(for EL, GM, Mx)

Week 8 - 07.04.25



Review of Fourier transform

Let $f: \mathbb{R} \rightarrow \mathbb{C}$ be a (piecewise) continuous function.



Fourier transform of f :

$$\hat{f}(\alpha) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-i\alpha x} f(x) dx$$

Sometimes we also use the notation

$$F[f](\alpha) \quad \text{for} \quad \hat{f}(\alpha).$$

In order for the integral to be well-defined:

We need $\int_{-\infty}^{+\infty} |f(x)| dx < \infty$ ($f \in L^1(\mathbb{R})$)

Inverse Fourier transform:

$$F^{-1}[f](x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{+i\alpha x} f(\alpha) d\alpha$$

It is the inverse of F :

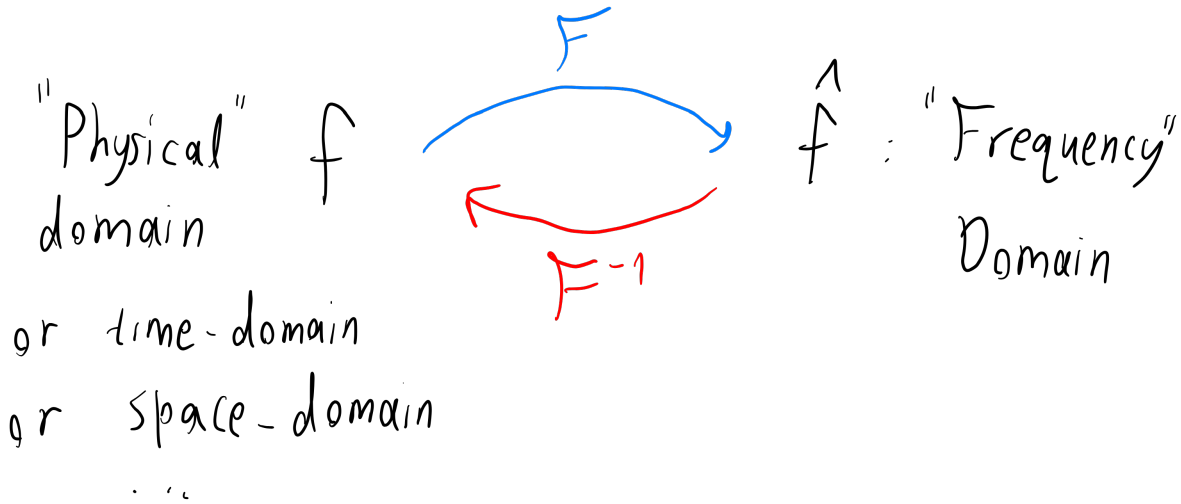
$$F^{-1}[F[f]] = f.$$

Note: The Fourier transform maps

$$\{f; \mathbb{R} \rightarrow \mathbb{C}\} \text{ to } \{\hat{f}; \mathbb{R} \rightarrow \mathbb{C}\}.$$

In general: Even if f is real valued, $\hat{f}$ might be $\mathbb{C}$ -valued.

(But, if f is real valued and even:
 $\hat{f}$ is also real valued)



Properties of the Fourier transform:

1) Linearity : $F[af + bg] = aF[f] + bF[g]$
if $a, b \in \mathbb{R}$

2) Time shift: Let $g(x) = f(x - x_0)$, then

$$\hat{g}(a) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-iax} g(x) dx$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-iax} f(x-x_0) dx$$

$$\stackrel{y=x-x_0}{=} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-ia(y+x_0)} f(y) dy$$

$$= \frac{1}{\sqrt{2\pi}} e^{-iax_0} \int_{-\infty}^{+\infty} e^{-ia y} f(y) dy$$

$$= e^{-iax_0} \hat{f}(a)$$

3) Frequency shift: Let $g(x) = e^{ia_0 x} f(x)$

Then:

$$\begin{aligned}
\hat{g}(\alpha) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-i\alpha x} g(x) dx \\
&= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-i\alpha x} \cdot e^{i\alpha_0 x} f(x) dx \\
&= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-i(\alpha - \alpha_0)x} f(x) dx \\
&= \hat{f}(\alpha - \alpha_0)
\end{aligned}$$

Note: Property ③ is essentially like property ② but for the inverse Fourier transform. This is to be expected, since the inverse Fourier transform is defined with a similar expression as the Fourier

transform (with e^{+iax} in place
of e^{-iax})

4) Re-scaling (or dilation)

Let $g(x) = f(\lambda x)$ for some $\lambda \neq 0$.

Then:

$$\hat{g}(a) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-iax} f(\lambda x) dx.$$

I want to substitute $y = \lambda x$. If

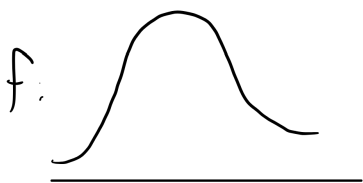
$\lambda > 0$, then I get $\int_{-\infty}^{+\infty} \dots \frac{dy}{\lambda}$,

while if $\lambda < 0$ I get $\int_{+\infty}^{-\infty} \dots \frac{dy}{\lambda} = \int_{-\infty}^{+\infty} \dots \frac{dy}{|\lambda|}$

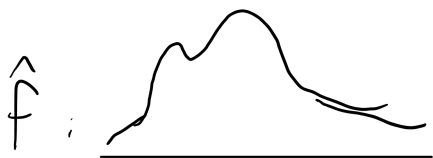
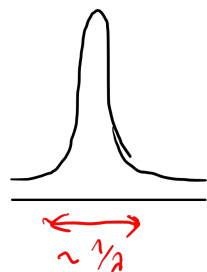
In both cases:

$$\begin{aligned}\hat{g}(\alpha) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-i\alpha \frac{y}{\lambda}} \cdot f(y) \frac{dy}{|\lambda|} \\ &= \frac{1}{|\lambda|} \cdot \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-i\frac{\alpha}{\lambda} y} f(y) dy \\ &= \frac{1}{|\lambda|} \hat{f}\left(\frac{\alpha}{\lambda}\right).\end{aligned}$$

Remark: If the graphs of f and $\hat{f}$ look as follows:



then, as $\lambda \rightarrow +\infty$,
the graph of
 $f(\lambda x)$ is



and of
 $\frac{1}{|\lambda|} \hat{f}\left(\frac{\alpha}{\lambda}\right)$:



So: When f is very concentrated around a point, $\mathcal{F}[f]$ is very spread out.

This is a manifestation of the uncertainty principle.

5) Derivatives and Fourier transform:

$$\text{Let } g(x) = f'(x).$$

$$\text{Then } \hat{g}(\alpha) = i\alpha \hat{f}(\alpha).$$

More generally:

$$\mathcal{F}[f^{(n)}](\alpha) = (i\alpha)^n \mathcal{F}[f](\alpha)$$

So the Fourier transform: Transforms differential equations with constant coefficients to algebraic equations on

the frequency domain.

Convolution: Given $f, g: \mathbb{R} \rightarrow \mathbb{C}$,
the convolution $f * g$ is defined by

$$f * g(x) = \int_{-\infty}^{+\infty} f(y) g(x-y) dy$$

Important properties:

i) Using the change of variable $\bar{y} = x - y$,
one can see that:

$$\int_{-\infty}^{+\infty} f(y) g(x-y) dy = \int_{-\infty}^{+\infty} f(x-\bar{y}) g(\bar{y}) d\bar{y}.$$

So:

$$f * g = g * f$$

ii)

$$f * (g * h) = (f * g) * h$$

$$\text{iii)} \quad f * (g+h) = f * g + f * h$$

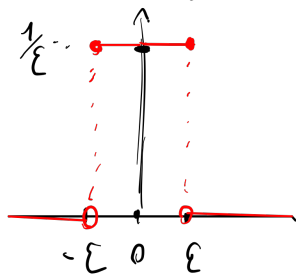
$$\text{iv)} \quad (f * g)' = f' * g = f * g'$$

So $f * g$ is at least as smooth as the smoother among f, g .

One way to interpret $f * g$: As the "local average" of f with weights g .

Important example:

$$\text{If } g(x) = \begin{cases} \frac{1}{2\varepsilon}, & x \in [-\varepsilon, \varepsilon], \\ 0, & \text{else,} \end{cases}$$



Then:

$$f * g(x) = \int_{-\infty}^{+\infty} f(y) g(x-y) dy =$$

$$= \frac{1}{2\varepsilon} \int_{x-\varepsilon}^{x+\varepsilon} f(y) dy$$

So it's really the local average of f ; Note that, in this case, as $\varepsilon \rightarrow 0$, $f * g(x) \rightarrow f(x)$.

The last point is true more generally:

If $g_\lambda(x)$ is a family of functions with $\int_{-\infty}^{+\infty} g_\lambda(x) dx = 1$ and $g_\lambda(x) = 0$ outside $[-\lambda, \lambda]$, then $f * g_\lambda(x) \xrightarrow{\lambda \rightarrow 0} f(x)$.

(In this case: As $\lambda \rightarrow 0$, $g_\lambda(x)$ looks like a δ -function).

6) Fourier transform and convolutions

Let $f, g: \mathbb{R} \rightarrow \mathbb{C}$, then

$$\mathcal{F}[f * g](\omega) = \sqrt{2\pi} \cdot \mathcal{F}[f](\omega) \cdot \mathcal{F}[g](\omega).$$

So: $\mathcal{F}$ transforms convolutions

into products (and the opposite as well, since $\mathcal{F}^{-1}$ has similar properties as $\mathcal{F}$).

In applications: We can often express

$\mathcal{F}[f]$ as the product of known functions,
in which case we recover f as
a convolution.

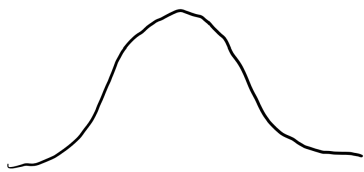
7) For the Gaussian function:

$$\text{If } f(x) = e^{-x^2/2}, \quad \hat{f}(a) = e^{-a^2/2}$$

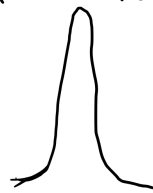
(We will prove it in the exercises).

So from property 4: For $\lambda \neq 0$,

$$\text{If } f_\lambda(x) = e^{-\lambda^2 \frac{x^2}{2}} \Rightarrow \hat{f}_\lambda(a) = \frac{1}{|\lambda|} e^{-\frac{a^2}{2\lambda^2}}$$



$\longleftrightarrow$
 $\frac{1}{\lambda}$



$\longleftrightarrow$
 λ

8) Plancherel's theorem:

$$\int_{-\infty}^{+\infty} |f(x)|^2 dx = \int_{-\infty}^{+\infty} |\hat{f}(\omega)|^2 d\omega$$

(same energy for a signal and its Fourier transform).

How does complex analysis enter the study of the Fourier transform?

Let's see an example:

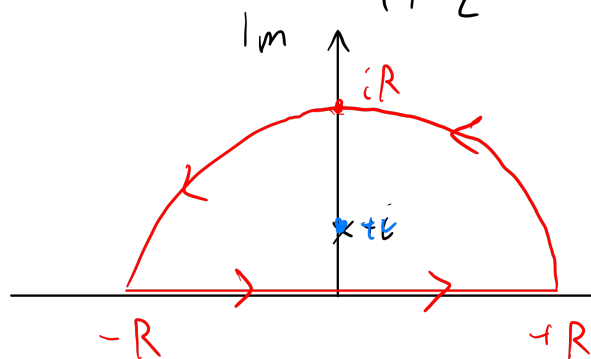
Let $f(x) = \frac{1}{1+x^2}$, then

$$\hat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-i\omega x} \frac{1}{1+x^2} dx.$$

We can evaluate this integral as an application of the residue theorem:

Let $a \leq 0$ (similar analysis if $a > 0$).

$$\text{If } f(z) = \frac{e^{-iaz}}{1+z^2} = \frac{e^{-iaz}}{(z-i)(z+i)}$$



Since $a \leq 0$:

On $\{|z|=R\} \cap \{\operatorname{Im}(z) > 0\}$
 f decays as $R \rightarrow +\infty$

From the examples we have already calculated in class (For $\frac{P(x)}{Q(x)} e^{iax}$, $a \geq 0$)

$$\begin{aligned} \int_{-\infty}^{+\infty} f(x) dx &= 2\pi i \operatorname{Res}_{+i}(f) = 2\pi i \lim_{z \rightarrow i} (z-i) f(z) \\ &= 2\pi i \cdot \frac{e^a}{2i} \end{aligned}$$

$$= \pi \cdot e^{-|a|}$$

So, for $a \leq 0$: $\mathcal{F}\left[\frac{1}{1+x^2}\right](a) = \frac{1}{\sqrt{2\pi}} \cdot \pi \cdot e^{-|a|}$

$$= \sqrt{\frac{\pi}{2}} e^{-|a|}$$

For $a \geq 0$: The same computation
 (but with the semi-circle now in
 the lower half plane) gives the
 same formula.

The Laplace transform:

The Fourier transform: Used for signals that are defined on the whole real line.

Laplace transform: For signals defined on $t \geq 0$.

Both: Transform differential equations into algebraic equations.

Let $f: [0, +\infty) \rightarrow \mathbb{C}$ be a piecewise continuous function. The Laplace transform of f :

$$\mathcal{L}[f](z) = \int_0^{+\infty} f(t) e^{-zt} dt$$

(z : complex!)

Sometimes: We use capital letters to denote the Laplace transform, for instance:

$$F(z) = \mathcal{L}[f](z), \quad G(z) = \mathcal{L}[g](z)$$

When is the integral well-defined (i.e. "converges")?

We need $\int_0^{+\infty} |f(t)| |e^{-zt}| dt < +\infty$.

Suppose that, for some $\gamma_0 \in \mathbb{R}$, we have

$$\int_0^{+\infty} |f(t)| e^{-\gamma_0 t} dt < +\infty$$

(naïvely, this condition says that,
as $t \rightarrow +\infty$, $|f(t)| < e^{\gamma_0 t}$...)

Then:

$$|\mathcal{L}[f](z)| \leq \int_0^{+\infty} |f(t)| |e^{-zt}| dt$$

$$= \int_0^{+\infty} |f(t)| \cdot |e^{-\operatorname{Re}(z) \cdot t - \operatorname{Im}(z) \cdot t \cdot i}| dt$$

Since

$$|e^{-\operatorname{Im}(z) \cdot t \cdot i}| = 1$$

$$= \int_0^{+\infty} |f(t)| e^{-\operatorname{Re}(z) \cdot t} dt$$

$$\stackrel{\operatorname{Re}(z) \geq \gamma_0}{\leq} \int_0^{+\infty} |f(t)| e^{-\gamma_0 t} dt$$

$$< +\infty.$$

Therefore, if such a γ_0 exists:

$\mathcal{L}[f](z)$ is defined at least over

$$U = \{z \in \mathbb{C} : \operatorname{Re}(z) > \gamma_0\}.$$

Definition:

A γ_0 such that $\mathcal{L}[f](z)$ is well-defined (as above) for any z with $\operatorname{Re}(z) > \gamma_0$ is called an **abscissa of convergence**.

(abscissa: Another name for the x -coordinate)

Example: If $f(t) = 0$ for $t \geq M$,

then
$$\int_0^{+\infty} |f(t)| e^{-\gamma_0 t} dt = \int_0^M |f(t)| e^{-\gamma_0 t} dt$$

$$< +\infty \quad \forall \gamma_0 \in \mathbb{R}$$

So in that case, $\mathcal{L}[f](z)$ is defined for $z \in \mathbb{C}$.

Remark: The "faster" the rate of growth for f as $t \rightarrow +\infty$ (or the "slower" the rate of decay), the smaller the domain $U \subseteq \mathbb{C}$ where $\mathcal{L}[f](z)$ is

well-defined.

If γ_0 is an abscissa of convergence
and $U = \{z : \operatorname{Re}(z) > \gamma_0\}$
then $\mathcal{L}[f]$ is holomorphic on U .

Let $F(z) = \mathcal{L}[f](z)$. Then

$$\begin{aligned}\frac{d}{dz} F(z) &= \frac{d}{dz} \int_0^{+\infty} f(t) e^{-zt} dt \\ &= \int_0^{+\infty} f(t) (-t) e^{-zt} dt \\ &= -\mathcal{L}[t \cdot f(t)](z).\end{aligned}$$

So:

Theorem: If $F(z) = \mathcal{L}[f](z)$,

$$F'(z) = -\mathcal{L}[t \cdot f(t)].$$

We will also use the following lemma (proven in the exercises):

Lemma: If γ_0 is an abscissa of convergence for $\mathcal{L}[f(t)]$, it is also for $\mathcal{L}[t \cdot f(t)]$.

Example 1.

Let $f: [0, +\infty) \rightarrow \mathbb{C}$ be $f(t) = 1$ (constant). Then, for any $\gamma_0 > 0$:

$$\begin{aligned}
 \int_0^{+\infty} |f(t)| e^{-\gamma_0 t} dt &= \int_0^{+\infty} e^{-\gamma_0 t} dt \\
 &= \left[\frac{e^{-\gamma_0 t}}{-\gamma_0} \right]_{t=0}^{t=+\infty} \\
 &= \frac{1}{\gamma_0} < \infty
 \end{aligned}$$

So any $\gamma_0 > 0$ is an abscissa of convergence $\Rightarrow \mathcal{L}[f](z)$ is defined (and holomorphic) for $\operatorname{Re}(z) > 0$.

And:

$$\begin{aligned}
 \mathcal{L}[f](z) &= \int_0^{+\infty} 1 \cdot e^{-z t} dt \quad z \neq 0 \\
 &= \left[\frac{e^{-z t}}{-z} \right]_{t=0}^{t=+\infty} =
 \end{aligned}$$

$$= -\frac{1}{z} \left(\lim_{t \rightarrow +\infty} (e^{-zt}) - e^{-0} \right)$$

But: $\lim_{t \rightarrow +\infty} (e^{-zt}) = 0$ for $\operatorname{Re}(z) > 0$,

since $|e^{-zt}| = |e^{-\operatorname{Re}(z)t}| \cdot |e^{-\operatorname{Im}(z)t \cdot i}|$



$\rightarrow 0$
as $t \rightarrow +\infty$



$= 1$

So: $\mathcal{L}[1](z) = \frac{1}{z}$

Example 2:

Let $f: [0, +\infty) \rightarrow \mathbb{C}$, $f(t) = t$.

$$\text{Then } \frac{d}{dz} \mathcal{L}[1](z) = -\mathcal{L}[t](z) \\ = -\mathcal{L}[f](z)$$

$$\text{So } \mathcal{L}[f](z) = -\frac{d}{dz} \left(\frac{1}{z} \right) = \frac{1}{z^2}$$

(Also holomorphic for $\operatorname{Re}(z) > 0$, since as we saw in the lemma, if γ_0 is an abscissa of convergence for $\mathcal{L}[f]$, the same is true for $\mathcal{L}[t \cdot f]$ and any $\gamma_0 > 0$ is an abscissa of convergence for $\mathcal{L}[1]$)

Repeating this construction:

Theorem: For any $n \in \mathbb{N}$,

$$\mathcal{L}[t^n](z) = \frac{n!}{z^{n+1}}$$

defined for $\operatorname{Re}(z) > 0$.